

ETW Branes and Inflationary Predictions for Rocky and Swampy Landscapes

Arthur Hebecker (Heidelberg)

based on work with **Bjoern Friedrich, Alexander Westphal and Sebastian Zell**

(cf. also earlier work with **Friedrich/Walcher/Salmhofer/Strauss**)

Outline

- The Measure Problem ist still there!
How is it affected by the 'Swampland Revolution' ?
- Our attempt to define an explicit, usable measure:
'Local Wheeler-DeWitt Measure'
- Input from 'Rocky Landscapes': KKLT, LVS etc.
- Input from 'Swampy Landscapes': maybe no (long-lived) dS,
Cobordism.
- Towards a prediction for the scale of inflation.

Introduction/Motivation

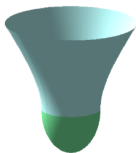
- **Generically:** Many vacua / Multiverse \Rightarrow Measure Problem.
Linde/Mezhlumian '93
- **Concretely:** Flux Landscape; “ 10^{500} vacua”
 \Rightarrow Measure Problem goes center stage.
Bousso/Polchinski '00, Denef/Douglas '04,
- With KKLT/LVS under pressure,
the (flux) landscape does not go away.
- Even if only slow-roll (or short-lived dS) exist...
....need a method to 'predict' our 'vacuum'

Introduction/Motivation (continued)

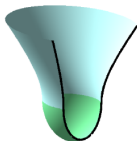
Key new 'swampland' input:

- Maybe no (multitude of) long-lived dS;
Maybe instead mostly/only slow-roll.
- Cobordism Conjecture \Rightarrow End-of-the-World Branes abundant
These ETW branes can be key players
in 'Creation from Nothing'.
(to be quantified below)

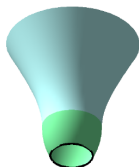
Preliminary illustration:



'No-Boundary'



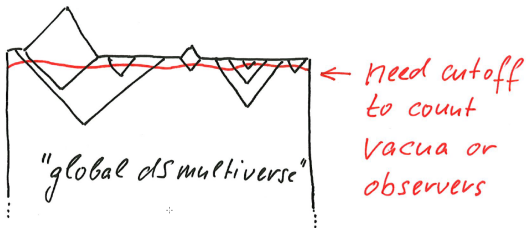
'Bubble-of-Something'



'Boundary proposal'

Measure problem and potentially decisive role of creation processes

- Standard view: Different vacua \rightarrow different patches in 'global dS multiverse'. Measure problem \equiv problem of cutoff choice.



- Based on the 'Cosmological Central Dogma', we want to argue for a more fundamental, quantum-mechanical measure.

Banks '01, Susskind '21

Friedrich/AH/Salmhofer/Strauss/Walcher '22,
Friedrich/AH/Westphal/Zell - to appear

Towards a 'Quantum-Measure'

- **Cosmological Central Dogma:**

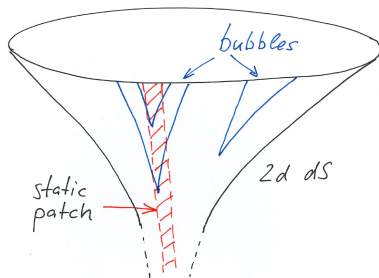
dS space is a finite system with $\dim(\mathcal{H}) = e^S$.

- Eternal Inflation \equiv Series of transitions between different subspaces (with $\dim(\mathcal{H}_i) = e^{S_i}$).

- Even better: Write down corresponding Wheeler-DeWitt equation:

$$H\psi = \chi$$

- Crucially, a source term for the creation from nothing is **unavoidable**.



'Local Wheeler-DeWitt Measure' (continued)

- Denote the sources by J_j and the decay rates by $\Gamma_{i \rightarrow j}$.
- Then the relevant rate equations read

$$J_i = \sum_{j \in dS} (p_j \Gamma_{i \rightarrow j} - p_i \Gamma_{j \rightarrow i}) + p_i \sum_{y \in Terminal} \Gamma_{i \rightarrow y} .$$

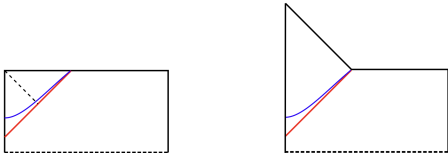
- The solution can be given as a series:

$$p_i = \frac{1}{\Gamma_i} \left\{ J_i + \sum_j J_j \frac{\Gamma_{j \rightarrow i}}{\Gamma_j} + \sum_{j,k} J_j \frac{\Gamma_{j \rightarrow k}}{\Gamma_j} \frac{\Gamma_{k \rightarrow i}}{\Gamma_k} + \dots \right\}$$

(Here Γ_i is the total decay rate of vacuum i .)

A conceptual problem: Reheating to Minkowski

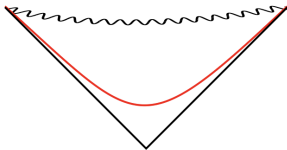
- As long as there are only dS and AdS vacua (and a non-zero rate for creation from nothing), finiteness is obvious.
- There is a sensitivity to the number of observers on the horizon-sized patch of the reheating surface.
But we ignore this (non-exponential!) effect.
- However, this changes once we include Minkowski-bubbles:



Now we have no reason to cut off the reheating surface at horizon size. **Technically, the projection $\|\psi|_i\|^2$ can be infinite.**

First Aside:

- One might think that this problem also arises for reheating in an AdS bubble. After all, $\dim(\mathcal{H}_{AdS}) = \infty$ and the reheating surface is infinitely large:



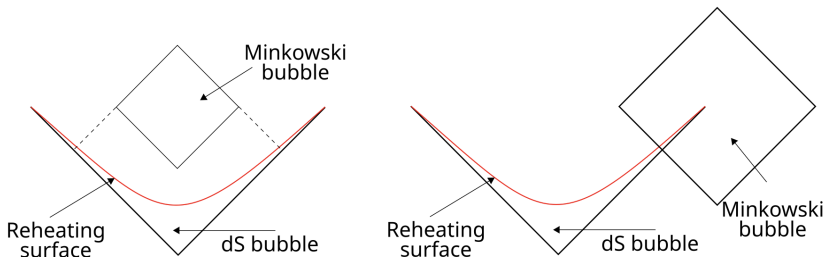
- However, we believe this can be dismissed because the future singularity ensures that there is **no infinity in any causally connected region.**

Second Aside:

- Maybe the problem is absent because there can be no observers on a Minkowski-space reheating surface (e.g. due to $\mathcal{N} = 2$ SUSY).

cf. Douglas '12

- However, even though Minkowski bubbles as such are in this case harmless, **bubble collisions are not!**

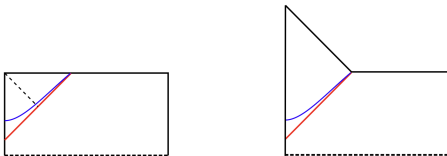


- What is worse: The observer-infinity in Minkowski depends on fine details of bubble-dynamics.

Kleban '11, Freivogel '11

Our proposal:

- Appeal to an ‘**Effective CCD**’, based on the **similarity** of the **reheating surfaces** in dS and Minkowski:



- In essence, we claim that even in Minkowski only a **finite portion** of the surface ($\sim 1/H_{\text{reh}}^3$) is independent.
- Finiteness is then regained even in the presence of bubbles with Minkowski-space reheating.

Alternative possibility:

- We could try to take the infinity of Minkowski-space reheating surfaces seriously (no redundancy).
- This would imply a key prediction: **The dark energy in our universe will decay** – our future is Minkowski space.

A Footnote:

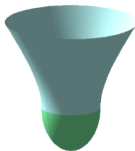
If no Minkowski-space reheating surfaces with observers exist in the landscape/multiverse, then **collision rates with Minkowski bubbles** determine the most likely vacuum.

... unsatisfactory....?

For now, we will use the **'Effective CCD'** logic....

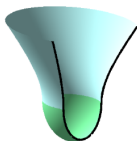
Towards explicit predictions

- We need **creation/decay rates**.
- In contrast to volume-weighted measures, our **local** measure crucially depends on **creation rates**. So let's start from those:



'No-Boundary'

Hartle/Hawking
Linde/Vilenkin



'Bubble-of-Something'

Hawking/Turok
Bousso/Chamblin
Garriga, Blanco-Pillado, ...

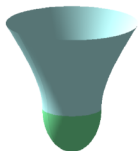


['Boundary proposal']

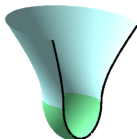
[Friedrich/AH]

[Cf. recent discussion of 'Bubble of Something' for String Landscape in Friedrich/AH/Walcher '23. Also, much recent work on inverse 'Bubble of Nothing' process: Garcia-Etxebarria/Montero/Sousa/Valenzuela, Draper et al., Angius/Calderon-Infante/Delgado/Huertas/Uranga, ...]

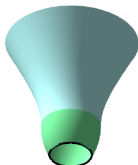
Creation Rates



'No-Boundary'

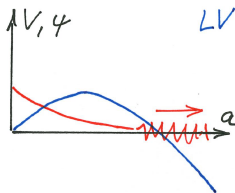
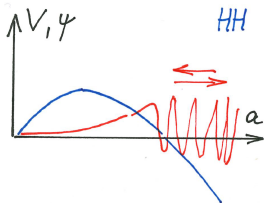


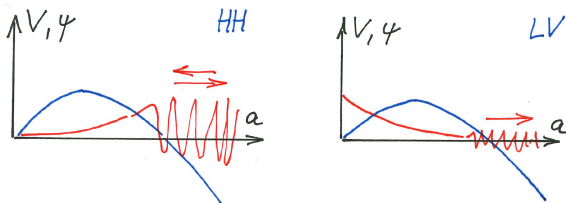
'Bubble-of-Something'



'Boundary proposal'

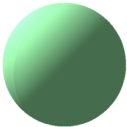
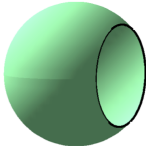
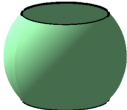
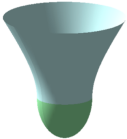
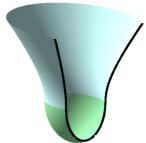
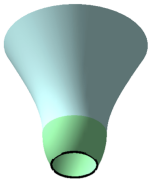
- A key question for all three processes is the sign in the exponent of the rate: $J \sim \exp(\pm S)$ ('LV vs. HH')
- Illustration of our (subjective, inconclusive) view:





- The (by definition real) HH version describes a 'ground state of the universe'. Maybe not suitable for 'creation rates'?
- Also, in strong tension with observation.
as recently quantified in Maldacena '24
- By contrast, the LV sign choice suffers from a 'matter-instability'. This may remove the exponential suppression.
Rubakov '84
- For the time being, we will remain open to both sign choices.

- Thus we have: $J \sim \exp(\pm S)$ with:

	No-Boundary (nb)	Bubble of Something (bos)	Boundary (b)
			
			
$S =$	$-8\pi^2 M_P^2 \ell_{dS}^2$	$-4\pi^2 M_P^2 \ell_{dS}^2 \left(1 - \frac{T\ell_{dS}}{\sqrt{T^2\ell_{dS}^2 + 4M_P^4}} \right)$	$-8\pi^2 M_P^2 \ell_{dS}^2 \sqrt{\frac{T^2\ell_{dS}^2}{T^2\ell_{dS}^2 + 4M_P^4}}$

\Rightarrow For LV, the 'bos'/'b' creation processes **always dominate** over 'nb' when the pos./neg.-tension ETW branes exist.

Another key concern:

- Small torus dS universes can expand from zero size **without any potential barrier**.
⇒ **no exponential suppression**.

Zeldovich/Starobinsky '84, Coule/Martin '99, Linde '04

- All dS vacua equally likely to be created (?)
- This **'creation with non-trivial topology'** deserves much more attention!

Next step toward predictions:

$$\text{Transition rates} \quad (\Gamma \sim \exp(-B))$$

Here only brief summary (see paper for more). We are building on KKLT/LVS-type flux vacua, but the conclusions look generic....

(1) Decay of the uplift / Decay by SUSY restoration:

$$B \sim T^4/(\Delta V)^2 \quad (\text{field theory regime, very fast})$$

(2) Decay to decompactification:

$$B \sim S_f - \mathcal{O}(1)S_f \quad (\text{much slower})$$

(3) Flux transitions:

$$B \sim S_f - M_P^6/T^2 \quad (\text{almost maximally suppressed})$$

Key conclusion: $\frac{\sum_{k \in \text{dS}} \Gamma_{j \rightarrow k}}{\Gamma_j} \ll 1$

(Transiting to any other dS is much less likely than terminal decay.)

- ⇒ Our solution-series converges fast.
- ⇒ May restrict attention to **direct creation from nothing** or **creation from nothing plus one tunnelling event**.
(i.e. only one or two step processes are relevant.)

Towards explicit predictions:

- Focus on observers on post-inflationary reheating surfaces (like us).
- Include inflationary plateaus as (short-lived) dS vacua '**inf(*i*)**', decaying to vacuum *i*.

⇒ Key formula:
$$p_{\text{inf}(i)} \simeq \frac{1}{\Gamma_{\text{inf}(i)}} \left(J_{\text{inf}(i)} + \sum_{o \neq \text{inf}(i)} J_o \frac{\Gamma_{o \rightarrow \text{inf}(i)}}{\Gamma_o} \right)$$

- **Question 1:** Does direct production (first term) or one-step tunnelling (second term) dominate?
- **Question 2:** What does this imply for the probability of 'observing' (in our past) a high- or low-scale inflationary plateau?
(for earlier analyses of this, cf. Pedro/Westphal '13)
- Our paper gives a detailed discussion of the answer, depending on various assumptions (see above....).
- Here, only one **'example answer'**:

Let's accept the LV sign, assume slow-roll vacua with high-tension ETW-branes exist ⇒ **Bubbles of something win!**
(Energy scale of Inflation determined by available ETW branes!)

Summary / Conclusions

- Predictions in any landscape (swampy or rocky) need a measure. I argued that, in a proper quantum approach, this is sensitive to 'Creation from Nothing'.
- Given the Cobordism Conjecture, a key ingredient in these creation events are ETW branes, allowing for 'BOS's or 'boundary processes'
- A key complication of the advertised 'local WDW measure' are Minkowski-space reheating surfaces.
(We suggested how to proceed, but this may not be final.)
- We derived simple equations giving explicit predictions, e.g. for the scale of inflation. But input is missing, including crucially :

ETW branes – tensions and availability?
Creation of small tori?